

$$r > R_3 \quad \varphi(r) = \frac{q}{4\pi\epsilon_0 r} = \frac{\sigma_3 R_3^2}{\epsilon_0 r}$$

$$r = R_3 \quad \varphi(R_3) = \frac{\sigma_3 R_3}{\epsilon_0}$$

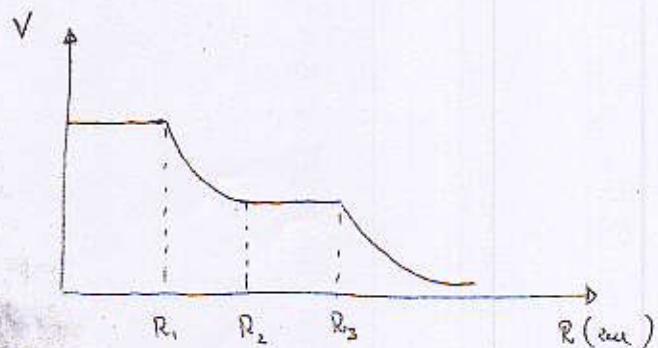
$$R_2 < r < R_3 \quad \varphi(R_2) = \varphi(R_3) = \varphi(r)$$

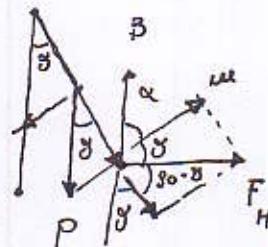
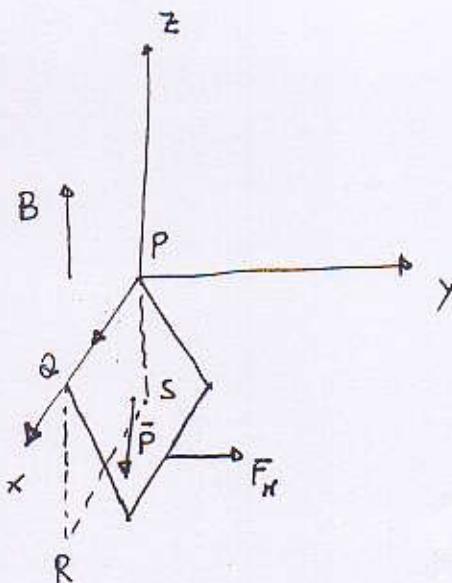
$$R_1 < r < R_2 \quad \Delta\varphi = - \int_{R_2}^r \vec{E} \cdot d\vec{l} = - \frac{1}{4\pi\epsilon_0} q \int_{R_2}^r \frac{dr}{r^2} = \\ = \frac{1}{4\pi\epsilon_0} \frac{q}{r} \Big|_{R_2}^r \Rightarrow \Delta\varphi = \frac{q}{4\pi\epsilon_0} \left(\frac{1}{r} - \frac{1}{R_2} \right)$$

$$\Rightarrow \varphi(r) - \varphi(R_3) = \frac{q}{4\pi\epsilon_0} \left(\frac{1}{r} - \frac{1}{R_2} \right) \Rightarrow \varphi(r) = \frac{q}{4\pi\epsilon_0} \left(\frac{1}{r} - \frac{1}{R_2} \right) + \varphi(R_3) =$$

$$= \frac{1}{\epsilon_0} \left(\frac{\sigma_1 R_1^2}{r} - \sigma_2 R_2 + \sigma_3 R_3 \right) \Rightarrow \varphi(R_1) = \frac{1}{\epsilon_0} (\sigma_1 R_1 + \sigma_2 R_2 + \sigma_3 R_3)$$

$$\Rightarrow \varphi(R_1) - \varphi(R_3) = \frac{1}{\epsilon_0} (\sigma_1 R_1 - \sigma_2 R_2)$$





$$F_{TP} = \mu_0 g \sin \gamma \Rightarrow H_p = \mu_0 g \frac{b}{2} \sin \gamma$$

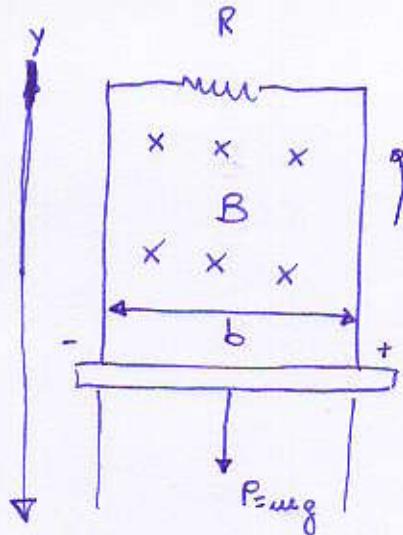
$$F_H = i a B \cos \gamma \Rightarrow H_H = i a b B \cos \gamma \sin(30^\circ)$$

$$\Rightarrow H_H = \bar{\mu} \times \bar{B} \quad \text{where } \bar{\mu} = i a b \cos \alpha$$

$$H_H = H_p \Rightarrow \mu_0 g \frac{b}{2} \sin \gamma = i a b B \cos \alpha$$

$$\mu_0 = 2(a+b)\delta \Rightarrow i = \frac{2(a+b)\delta g}{a B} \bar{T}_g \delta = 2.12 A$$

$$W = \int_0^{30^\circ} H_H d\gamma = i a b B \int_0^{30^\circ} \cos \gamma d\gamma = 4.27 \times 10^{-4} J$$



Il campo magnetico, ortogonale al piano, può avere seletti verso e uscita; si verifica, infatti, che il risultato è sempre una forza magnetica che agisce in opposizione alla forza peso.

Considero il sistema di ref. indicato in fig.

$$\mathcal{E} = - \frac{d\phi(B)}{dt}$$

$$d\phi(B) = -B b v dt \Rightarrow \mathcal{E} = + B b v \Rightarrow i = \frac{B b v}{R}$$

$$F_m = +i \vec{b} \times \vec{B} = -\frac{B^2 b^2}{R} v \quad \text{pulsione negativa}$$

$$-F_m + P = m \omega \Rightarrow \frac{B^2 b^2}{R} v(t) - m g = m \frac{dv(t)}{dt}$$

$$\frac{dv(t)}{dt} + \frac{B^2 b^2}{mR} v(t) = +g \quad , \quad \kappa = \frac{B^2 b^2}{mR}$$

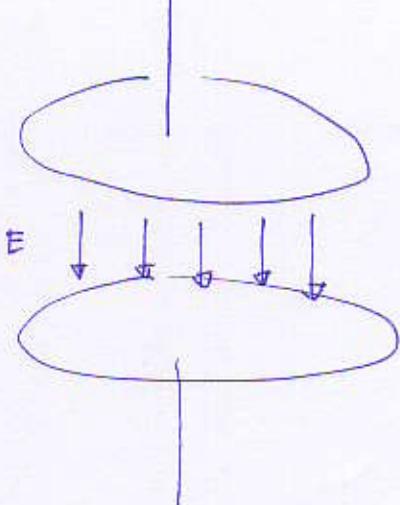
$$v(t) + \kappa v(t) = +g \Rightarrow \text{stabilità assoluta}$$

$$\text{O.M.G.: } v(t) = v_0 e^{-kt}$$

$$\text{SOL. PART: } V_p = \frac{g}{k}$$

$$\Rightarrow V_{TOT}(t) = \frac{g}{k} (1 - e^{-kt}) \Rightarrow V_0 = \frac{g}{k} \quad i_0 = \frac{B b g}{k R}$$

$$i(t) = \frac{B b g}{R k} (1 - e^{-kt})$$



$$\int_{\Gamma} \bar{B} \cdot d\bar{l} = \mu_0 \epsilon_0 \int \frac{\partial E}{\partial t} \cdot \bar{v}_n dS = \\ = \mu_0 \epsilon_0 \frac{\partial \phi(E)}{\partial t}$$

$\int_{\Gamma} \bar{B} \cdot d\bar{l} = 2\pi c B$ si e' scelto un percorso chiuso di forma circolare e di raggio r .

$$\int \frac{\partial E}{\partial t} \cdot \bar{v}_n dS = \pi r^2 \frac{\partial E}{\partial t} \Rightarrow 2\pi c B = \epsilon_0 \mu_0 c^2 E_0 \omega \cos \omega t$$

$$\Rightarrow B(t) = \frac{c \omega E_0}{2c^2} \cos \omega t$$