

A closer look at Critical Points

Slow dynamics, aging and their universal features

Andrea Gambassi



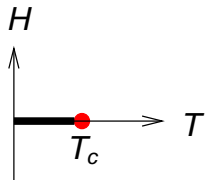
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Stuttgart



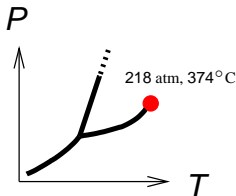
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Florian Baumann (Erlangen→Nancy)

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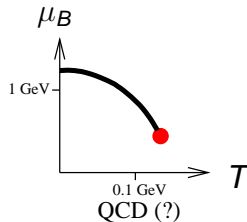
Critical Points & Critical Dynamics (reminder)



Ising



water



QCD (?)

Order parameter: $\phi_x(t)$

Correlated fluctuations: ξ, T_R

Hallmarks: ST $\xi \gg r_{\text{micr}}$ *collective*

DY $T_R \gg \tau_{\text{micr}}$ *slow*

\Rightarrow **Universality**:
 » experimental fact: $\chi \sim a \xi^{-\gamma/\nu}$
 » **minimal models** (effective theories, FT, RG and all that)

Critical Dynamics

TH:

late '70s: Hohenberg, Halperin...
basic minimal models

...theoretical work...

2005: Folk, Moser (review)

EX:

'70-'80s: neutron scattering
 $S(\mathbf{q}, \omega)$, magn.

...exp. difficulties...

2005: X-ray^a
“watching” critical relax.

^aMocuta et al. Science **308**, 1287

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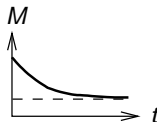
Mostly on: (a) Equilibrium DY close to CP

» $S(\mathbf{q}, \omega)$

» linear response κ, η , etc. $T_R^{(l)} \sim \xi^z$

(b) Non-linear relaxation

» $M(t) = \langle \phi_x(t) \rangle$, $T_R^{(nl)} \sim \xi^{z-\beta/\nu}$



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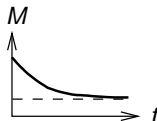
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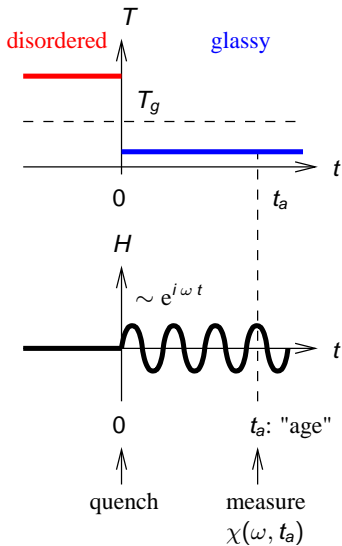
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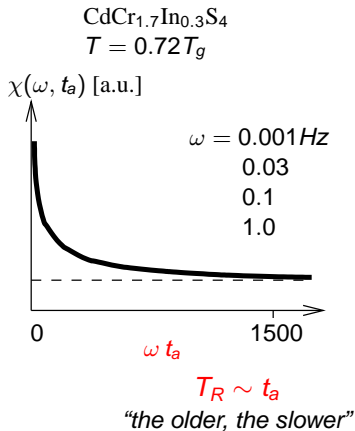
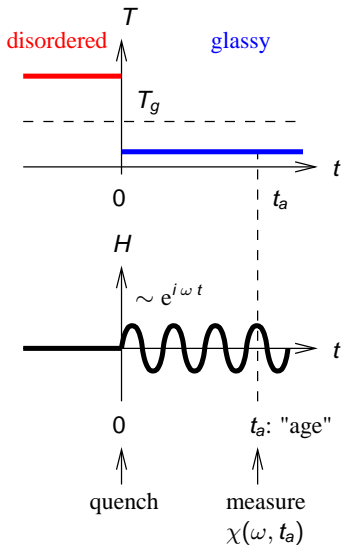


Two-time quantities display aging!_[CKP'94]

Aging (spin glasses)



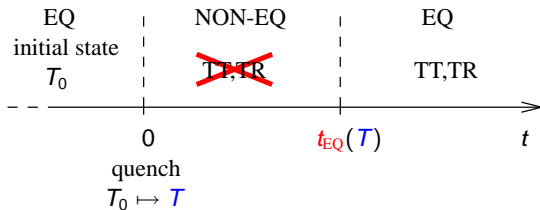
Aging (spin glasses)



[see Vincent et al, cond-mat/9607224]

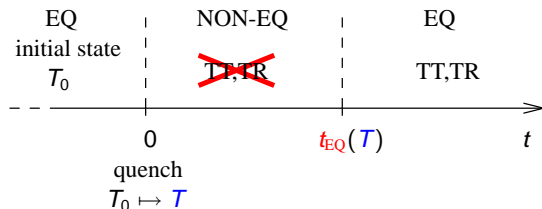
Relaxation

Ferromagnet [$T_c, \phi_x(t)$] relaxing towards *equilibrium*



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$$T > T_c : t_{EQ} < \infty \implies \text{EQ}$$

$$T < T_c : t_{EQ} = \infty \implies \text{phase ordering dynamics}$$

$$\boxed{T = T_c} : t_{EQ} = \infty \implies \text{Critical Dynamics}$$
$$\xi(t) \sim t^{1/z}$$

Dynamic Observables I

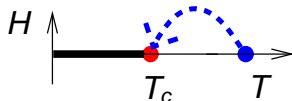
$\phi_x(t)$ order param. (eg, *local* fluct. magn.)

Simplest obs.: (1) one-time quantities M
(2) two- R, C

$$M(t) \equiv V^{-1} \sum_x \langle \phi_x(t) \rangle \quad \text{global} [= \langle \phi_{q=0}(t) \rangle]$$

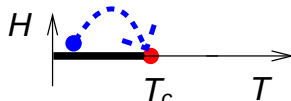
$$\boxed{T = T_c} :$$

$$M(0) = 0$$



$$\Rightarrow M(t) = 0, \quad \forall t$$

$$M(0) \neq 0$$



$$\Rightarrow M(t \rightarrow \infty) \sim t^{-\beta/\nu z}$$

$$M(t \rightarrow \infty) = 0 \Rightarrow \text{no infos on dynamics!}$$

Dynamic Observables II

$$s < t \quad \begin{cases} C_{x-y}(t, s) \equiv \langle \phi_x(t) \phi_y(s) \rangle_{\text{conn}} & \text{CORR} \\ R_{x-y}(t, s) \equiv \frac{\delta}{\delta h_y(s)} \langle \phi_x(t) \rangle \Big|_{h=0} & \text{RESP (lin.)} \end{cases}$$

$$t_{\text{EQ}}(T) \ll s < t: \quad C_x(t, s) = C_x^{(\text{eq})}(t-s) \quad ; \quad R_x(t, s) = R_x^{(\text{eq})}(t-s)$$

$$\boxed{TR_x^{(\text{eq})}(\tau) = -\frac{dC_x^{(\text{eq})}(\tau)}{d\tau}} \quad \text{FDT}$$

In general? **Fluctuation-Dissipation Ratio** [CK '94]

$$X(t, s) \equiv \frac{TR(t, s)}{\partial_s C(t, s)}$$

$$X^\infty \equiv \lim_{s \rightarrow \infty} \lim_{t \rightarrow \infty} X(t, s) \quad = 1 \text{ if } t_{\text{EQ}}(T) < \infty$$

??? FDT & TD with $T_{\text{eff}} \equiv T/X^\infty$ $\overset{\text{YES}}{\rightsquigarrow}$ ∞ -range glass

Model (Glauber dyn)		$X^\infty (\infty \mapsto T_c)$
Random Walk, GF	[1]	1/2
Spherical	[2]	$1 - 2/d$
1-dim. Ising	[2,3]	1/2
2-dim. Ising	[2]	0.26(1)
“ “	[4]	0.340(5)
“ “	[5]	0.33(2)
“ “	[6]	0.33(1)
“ “	[7]	0.330(5)
3-dim. Ising	[2]	0.40
3-dim. XY	[8]	0.43(4)

Exact solution, Monte Carlo simulations.

[1] Cugliandolo, Kurchan, Parisi 1994

[2] Godrèche, Luck 1999, 2000

[3] Zannetti *et al.* 1999

[4] Mayer *et al.* 2003

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Idea: X^∞ at $T = T_c$ is *universal*
[Godrèche, Luck '00]

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exploit Universality! [CG'02]

LATTICE, \mathbb{Z}^d

CONTINUUM, \mathbb{R}^d

Ising, $S_i(t)$
 $O(n)$, S_i^α

Scaling
↦

LGW, $\varphi(x, t)$
 $O(n)$ LGW, φ^α

1. spin-flip

Scaling
↦

MODEL A

2. spin-exch

Scaling
↦

MODEL B



analytical predictions for
 $X(t, s)$, X^∞ , scaling forms, exponents, etc.

Model A

one-component o.p. $\varphi(\mathbf{x}, t)$ on the continuum

$$\partial_t \varphi(\mathbf{x}, t) = -D \frac{\delta \mathcal{H}[\varphi]}{\delta \varphi(\mathbf{x}, t)} + \zeta(\mathbf{x}, t)$$

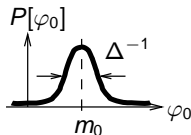
$$\langle \zeta(\mathbf{x}, t) \zeta(\mathbf{x}', t') \rangle = 2D \delta(\mathbf{x} - \mathbf{x}') \delta(t - t')$$

$$\Rightarrow P_{\text{EQ}}[\varphi] \propto e^{-\mathcal{H}[\varphi]}$$

$$\mathcal{H}[\varphi] = \int d^d \mathbf{x} \left[\frac{1}{2} (\nabla \varphi)^2 + \frac{1}{2} r \varphi^2 + \frac{u}{4!} \varphi^4 \right]$$

$$\text{[MSR'73, BJW'76]} \Rightarrow S[\varphi, \tilde{\varphi}] = \int_0^\infty dt \int d^d \mathbf{x} \left[\tilde{\varphi} \partial_t \varphi + D \tilde{\varphi} \frac{\delta \mathcal{H}}{\delta \varphi} - D \tilde{\varphi}^2 \right]$$

$$\langle \mathcal{O} \rangle_\zeta = \int [d\varphi d\tilde{\varphi}] \mathcal{O} e^{-S[\varphi, \tilde{\varphi}]} \quad \tilde{\varphi} \leftrightarrow h$$



initial cond. $\varphi_0(\mathbf{x}, t) \equiv \varphi(\mathbf{x}, t = 0)$ with

$$P[\varphi_0] \propto e^{-H_0[\varphi_0]}$$

$$H_0[\varphi_0] = \int d^d \mathbf{x} \frac{\Delta}{2} [\varphi_0(\mathbf{x}) - m_0]^2$$

[JSS'88] \Rightarrow FT with action $S[\varphi, \tilde{\varphi}] + H_0[\varphi_0]$

Scaling forms

$$T = T_C, \quad s < t \quad \left\{ \begin{array}{l} R_{q=0}(t, s) = A_R (t - s)^a \left(\frac{t}{s}\right)^\theta \mathcal{F}_R(s/t, t/t_0) \\ C_{q=0}(t, s) = A_C s (t - s)^a \left(\frac{t}{s}\right)^\theta \mathcal{F}_C(\quad, \quad) \end{array} \right.$$

$$a = (2 - \eta - z)/z$$

θ initial-slip exponent [JSS'88]

$\mathcal{F}_{R,C}(0, 0) = 1$, universal;

$$t_0 \equiv A_m m_0^{-1/\sigma} \text{ non-univ}$$

$$\sigma = \theta + a + \beta/(\nu z)$$

A_R, A_C non-univ.

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$$m_0 = \infty \quad \text{ie, } s, t \gg t_0 \quad [\text{CGK'06}]$$

$$R_{q=0} = a_R (t - s)^a \left(\frac{s}{t}\right)^{\beta\delta/(\nu z)} f_R(s/t)$$

$$\stackrel{t \gg s}{\approx} a_R t^a (s/t)^{\beta\delta/(\nu z)}$$

$$X^\infty = \frac{a_R}{a_c(1 - \frac{\beta\delta}{\nu z})}$$

\Rightarrow for long times: (a) $m_0 = 0$ (b) $m_0 \neq 0 \rightsquigarrow$ crossover

Results $m_0 = 0$

Ising model, Glauber dynamics

- **Global magnetization** $M(t)$:

X^∞	FT
$d = 2$	0.30(5)
$d = 3$	0.429(6)
$d > 4$	1/2

^aMayer, Berthier, Garrahan and Sollich (2003)...

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note: same results for
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 $\langle \varphi(\mathbf{x}, t) \rangle$

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- Other observables:

$R, C, X \mapsto R_{\mathcal{O}}, C_{\mathcal{O}}, X_{\mathcal{O}}$ (eg. $\mathcal{O} = \varphi^2$ energy density, $h_{\mathcal{O}}$ bath temp.)

Gaussian fluct. $(d > 4): X^\infty = X_{\mathcal{O}}^\infty \quad \forall \mathcal{O} \text{ local}$

Non-Gaussian fluct. $(d < 4): X^\infty \neq X_{\mathcal{O}}^\infty \implies ??? T_{\text{eff}} \equiv T_c / X^\infty$
!! MC $X_E^\infty = 0.33(2)$ ^(b)

- **Global magnetization** $M(t)$:

(a) Checked scaling forms & exp. $\beta\delta/(vz)$, MC $d = 2$

X^∞	FT	
$d = 2$	0.75	(^a)
$d > 4$	4/5	

note: same results for
local magnetization
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^aCalabrese, G, and Krzakala (2006)

- **Global magnetization** $M(t)$:

(a) Checked scaling forms & exp. $\beta\delta/(vz)$, MC $d = 2$

X^∞	FT	MC
(b) $d = 2$	0.75	0.73(1) ^(a)
$d > 4$	4/5	

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- $O(n)$, $n > 1$:** *unphysical dyn.*

Gaussian fluc.: $X_L^\infty = 4/5$
 $X_T^\infty = 2/3 \implies$ no T_{eff} even at the Gaussian level

- Purely dissipative dyn (Model A), $O(n)$ GLW:
 - $n = 1$ *anisotropic magnets/alloys*
 - $\forall n$ lattice spin models with **Glauber** dyn
 - ↪ $X^\infty(m_0 = 0)[2L]$: **OK** w. MC Ising 2d & 3d and XY 3d
 - ↪ $X^\infty(m_0 \neq 0)[1L]$: **OK** w. MC Ising 2d
 - ↪ $T_{\text{eff}} ???$
- Conserved dyn (Model B) scalar GLW
 - some *uniaxial ferromagnets*
 - Ising model with **Kawasaki** dyn
 - ↪ $[1L] X_{\text{Model B}}^\infty > X_{\text{Model A}}^\infty$: **OK** with MC Ising 2d [GKRT'04]
- Other dyn: Model C [1L], Model A dilute Ising [1L] [SP'04]

... what should be looked at...

? Aging at surfaces (different Universality Classes!) [MC:P'05; Sph mod:PB'05]

? More realistic dynamics...

- magnets \implies Models J,G (A)
- fluids \implies Model H (B)
... **waiting for experiments!**

? Slow dynamics at QCP

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- ? More realistic dynamics...
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... waiting for experiments!
- ? Slow dynamics at QCP

Take-home message:

FT is a viable approach to investigate aging phenomena at critical points!

Ref.: Calabrese&G, J.Phys.A **38** (2005) R133-R193