

Fluctuations in quenched spin-glasses and coarsening systems: are they similar ?

Federico Corberi

University of Salerno, Italy

Leticia F. Cugliandolo, Claudio Chamon

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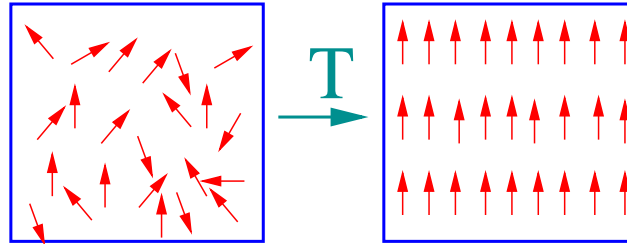
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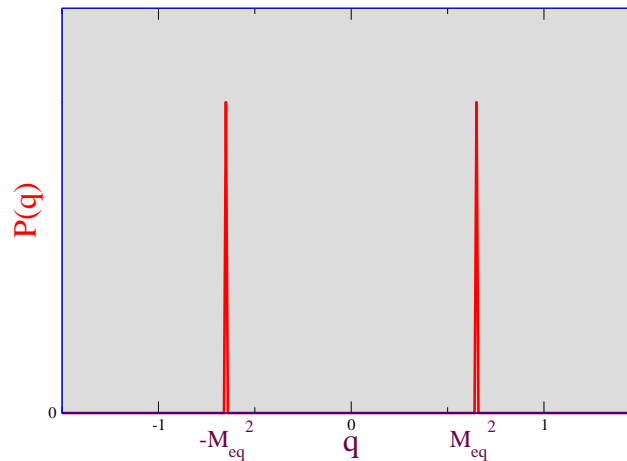
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- TRI: holding in glasses but not in magnets
- Numerical results
- Conclusions

Equilibrium - Statics

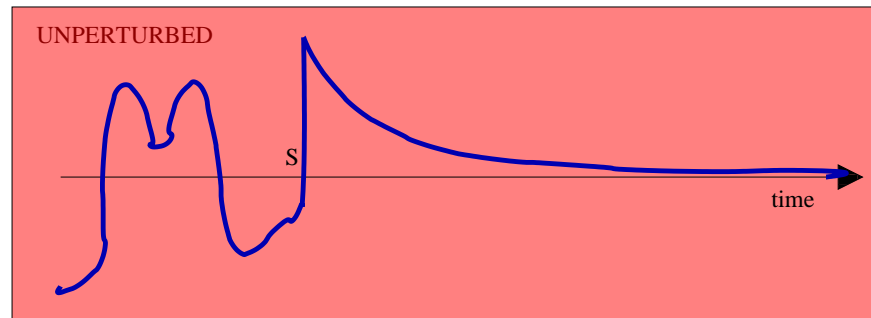
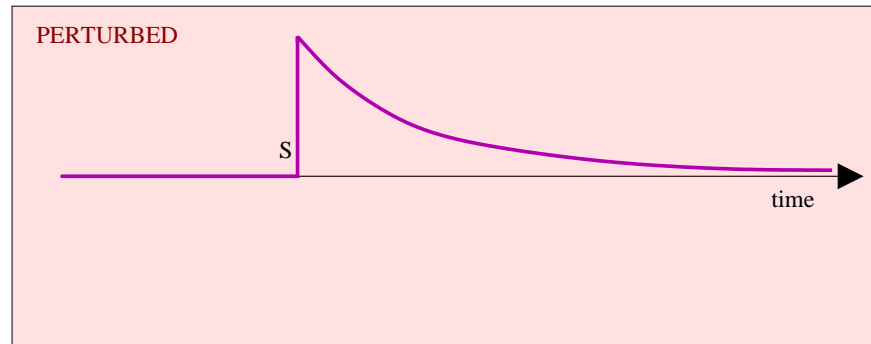


$$q^{\alpha\beta} = (1/N) \sum_{i=1}^N \sigma_i^\alpha \sigma_i^\beta \quad \text{Overlap}$$



Quasi-Equilibrium Dynamics

$$R_i(t, s) = \frac{\partial \langle \sigma_i(t) \rangle}{\partial h_i(s)} \quad h \rightarrow 0$$



Quasi-Equilibrium Dynamics: FDT



$$C_i(t, s) = \langle \sigma_i(t) \sigma_i(s) \rangle - \langle \sigma_i(t) \rangle \langle \sigma_i(s) \rangle$$

Quasi-Equilibrium Dynamics: FDT



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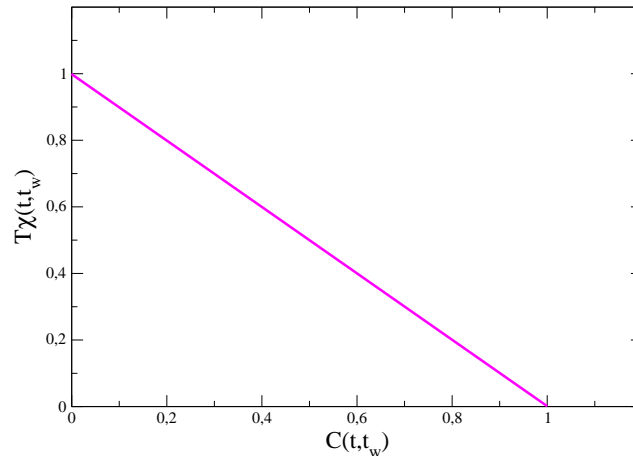
$$C_i(t, s) = C_i(t - s); \quad R_i(t, s) = R_i(t - s)$$

Quasi-Equilibrium Dynamics: FDT

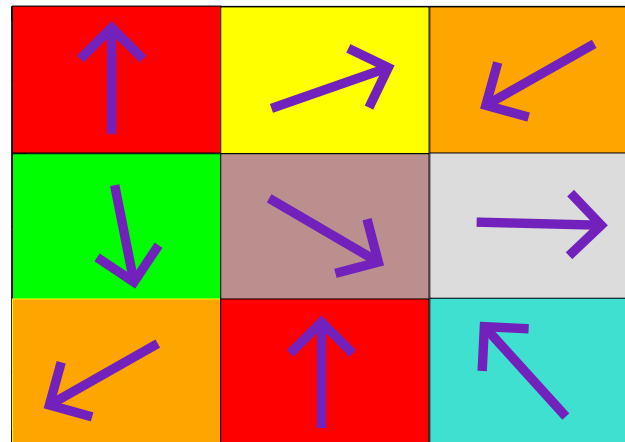
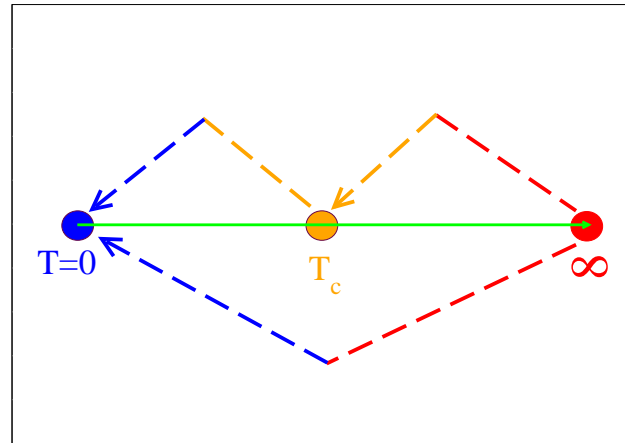
- $$C_i(t, s) = \langle \sigma_i(t) \sigma_i(s) \rangle - \langle \sigma_i(t) \rangle \langle \sigma_i(s) \rangle$$

- $$C_i(t, s) = C_i(t - s); \quad R_i(t, s) = R_i(t - s)$$

- $$\chi_i = f(C_i); \quad f \text{ is linear}; \quad \chi_i(t, s) = \int_s^t R_i(t, z) dz$$



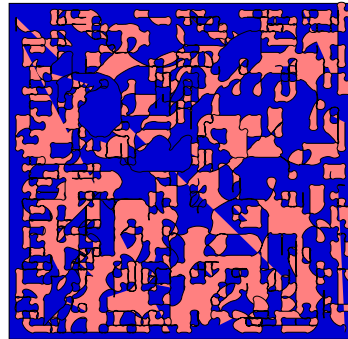
Far from Equilibrium: Quenches



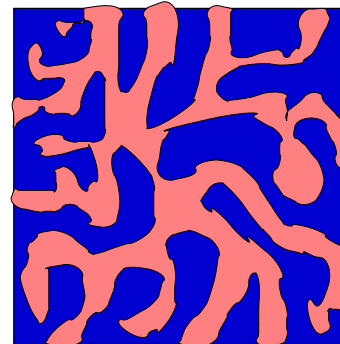
Far from Equilibrium-Growth of order

Size of *domains* $L(t)$ is growing in time.

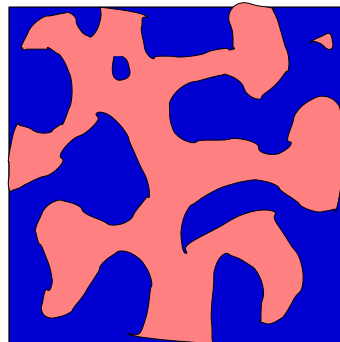
$$L(t) \simeq t^{1/z} \quad \textit{usually}$$



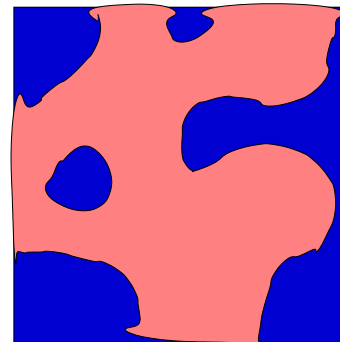
t_1



t_2



t_3



t_4

Far from Equilibrium - Scaling

$$C_i(t, s) = L(s)^{bz} f_C[L(t)/L(s)]$$

$$\chi_i(t, s) = L(s)^{az} f_\chi[L(t)/L(s)]$$

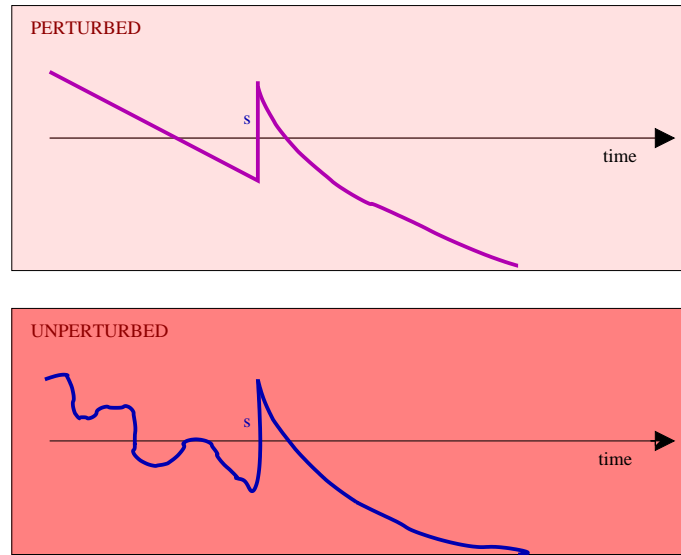
Analogous to (static) equilibrium scaling, i.e.

$$g(r, T) = \xi(T)^{-(d-2+\eta)} f_g[r/\xi(T)]$$

Using $L(t) \simeq t^{1/z} \Rightarrow$

$$C_i(t, s) = s^b \tilde{f}_C[t/s]; \quad \chi_i(t, s) = s^a \tilde{f}_\chi[t/s]$$

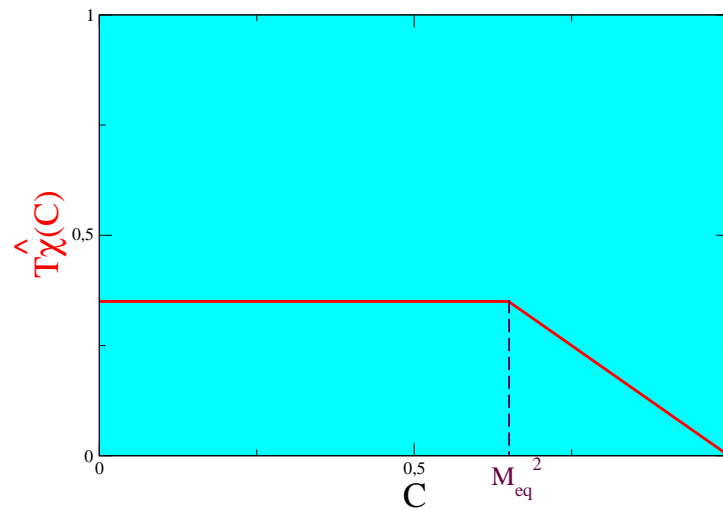
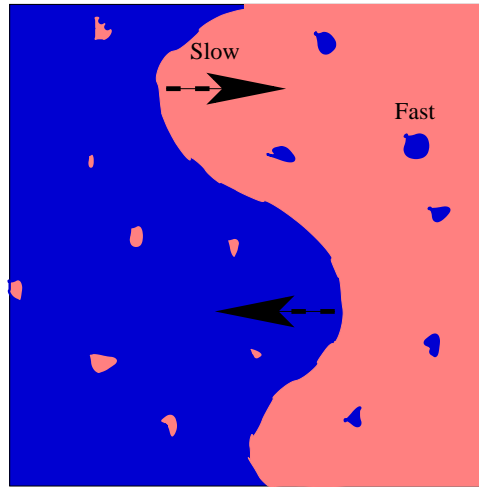
Far-from-Equilibrium: FDT



$$\chi_i = g(C_i, D_i); \quad D_i \text{ is another correlation function}^a$$

^a L.F. Cugliandolo, J. Kurchan and G. Parisi, J.Phys.I France 4, 1641 (1994);
E.Lippiello, F.Corberi, M.Zannetti, Phys. Rev. E 71, 036104 (2005); M.Baiesi,
C.Maes, and B.Wynants, Phys. Rev. Lett. 103, 010602 (2009).

Far-from-Equilibrium-2 components



Classification of systems (T_{eff})

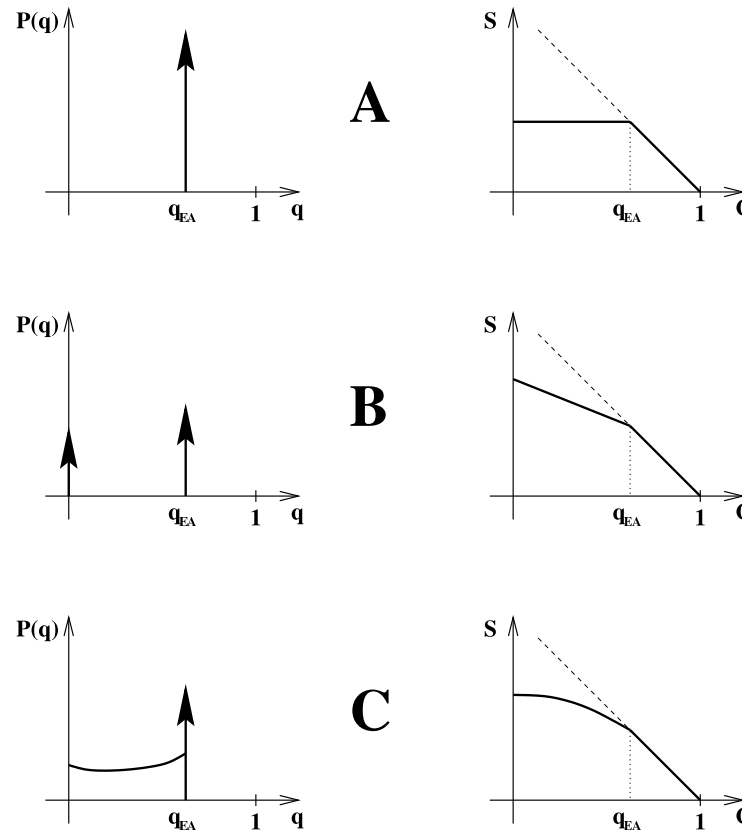
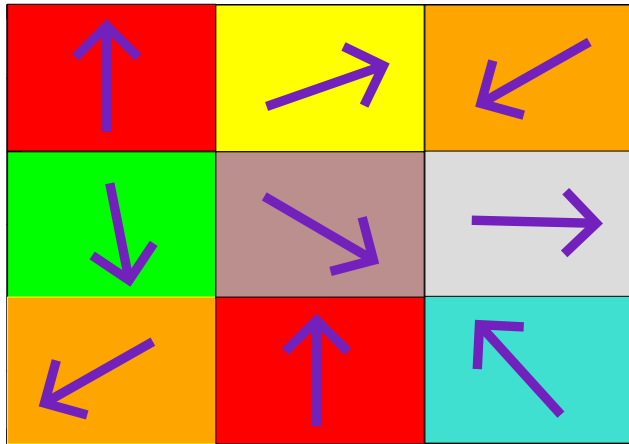


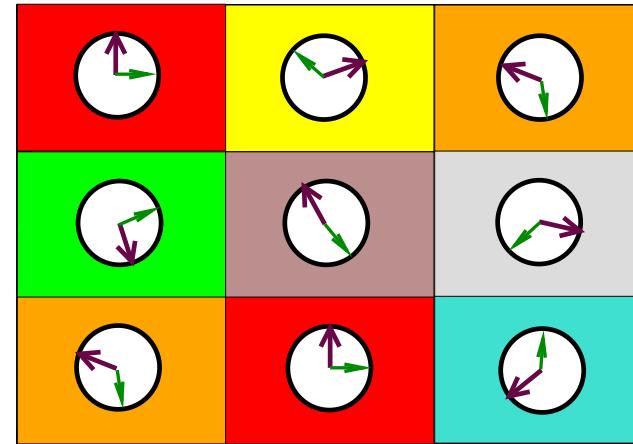
Figure 10: A possible model classification based on the function $S(C)$. The big arrows represent delta functions.

Where are finite dimensional SG and real glasses ?

t-reparametrization invariance-TRI



$O(N)$ symmetry



$t \rightarrow h(t); C \rightarrow C'$

TRI should hold true for glassy systems with $T_{eff} < \infty$ but not for coarsening with $T_{eff} = \infty$.

Fluctuating parts

$$\hat{C}_i(t, s) = [\sigma_i(t) - \langle \sigma_i(t) \rangle][\sigma_i(s) - \langle \sigma_i(s) \rangle]$$

$$\text{or } \hat{C}_i(t, s) = \sum_{\text{box around } i} \hat{C}_i(t, s)$$

Define properly also $\hat{\chi}_i(t, s)$.

$$V_{ij}^{CC}(t, s) = \langle \hat{C}_i(t, s) \hat{C}_j(t, s) \rangle - \langle \hat{C}_i(t, s) \rangle \langle \hat{C}_j(t, s) \rangle$$

$$V_{k=0}^{CC}(t, s) = \sum_{ij} V_{ij}^{CC}(t, s)$$

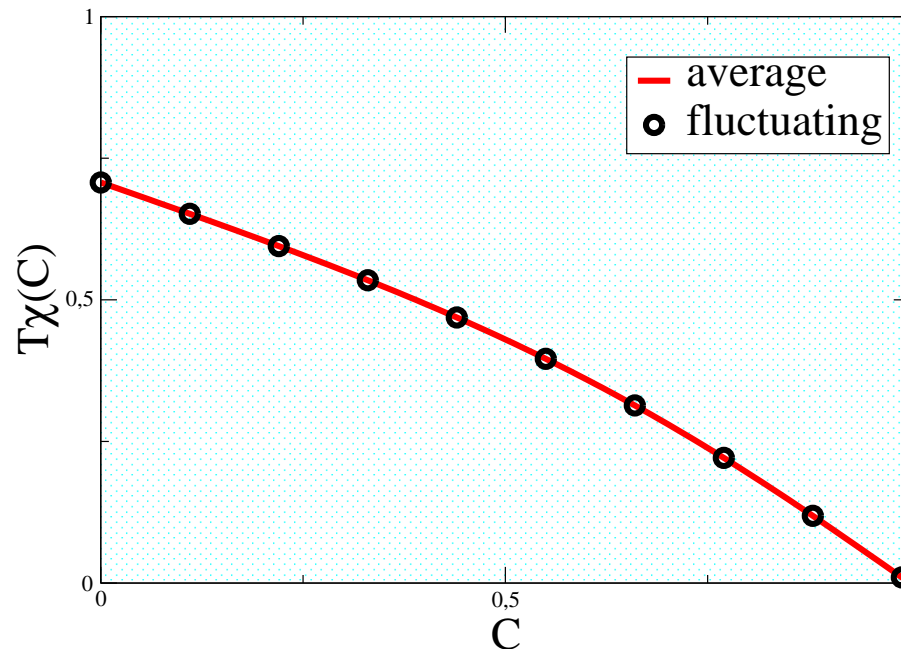
and similarly for V^{XX} and V^{CX} .

Generic variance is denoted V^{XY} onwards.

Transformations of $\hat{C}_i, \hat{\chi}_i$ under TRI

In a glassy system with $T_{eff} < \infty$, \hat{C}_i and $\hat{\chi}_i$ transform *in the same way* when $t \rightarrow h(t)$.

Physical meaning: all sort of fluctuations of $\hat{\chi}_i$ are accounted for by the fluctuations of the local clock \hat{C}_i .



Transformations of V^{XY} under TRI

If TRI holds also all the moments of \hat{C} and $\hat{\chi}$ must transform *in the same way*.

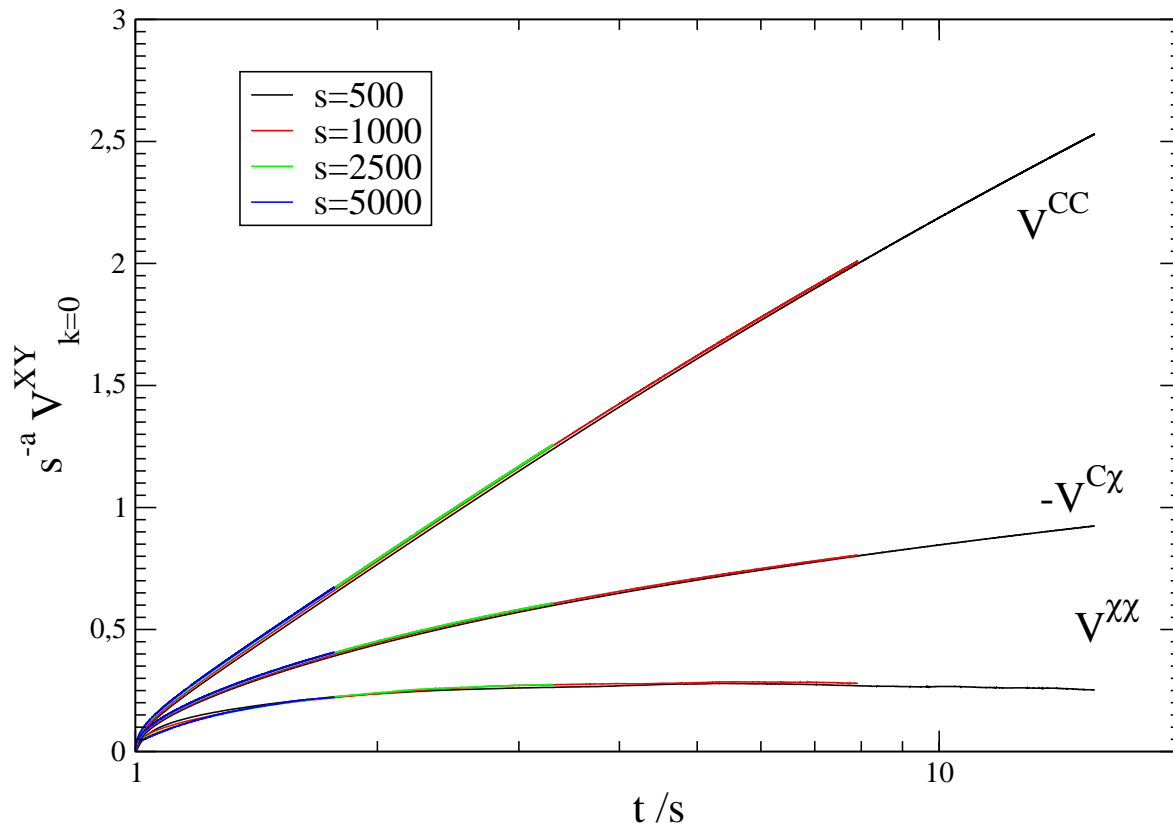
If scaling is obeyed in the form:

$$V^{XY}(t, s) = s^{a_{XY}} f_{XY}[t/s]$$

one must have the same exponent and behavior of f_{XY} .

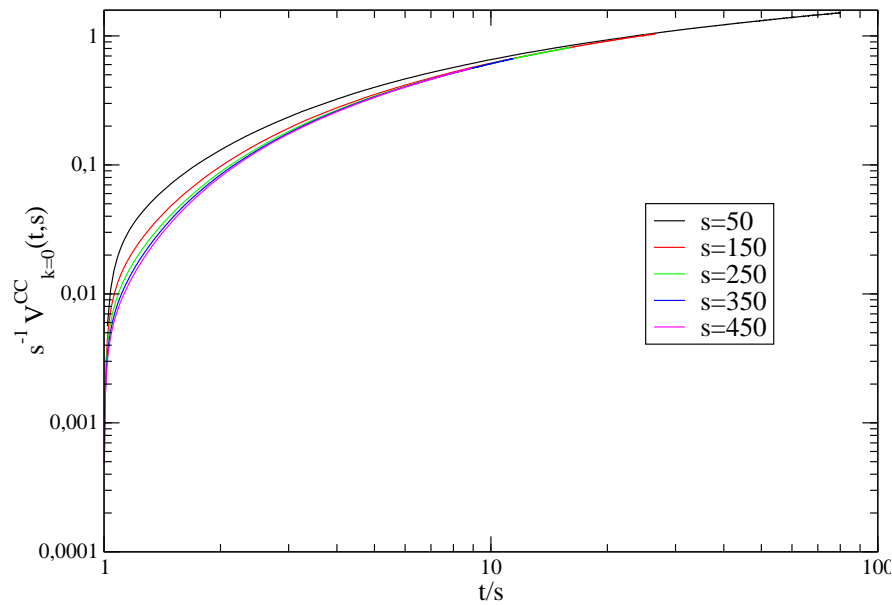
Strong prediction !

Numerical results SG 3d

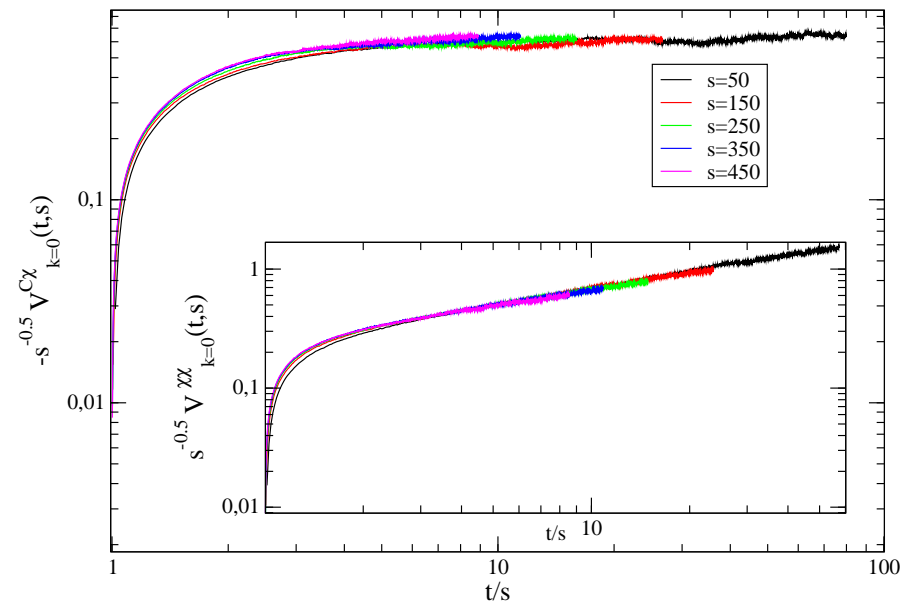


$a = 0.3$ for all V^{XY}

Numerical results Ising 2d

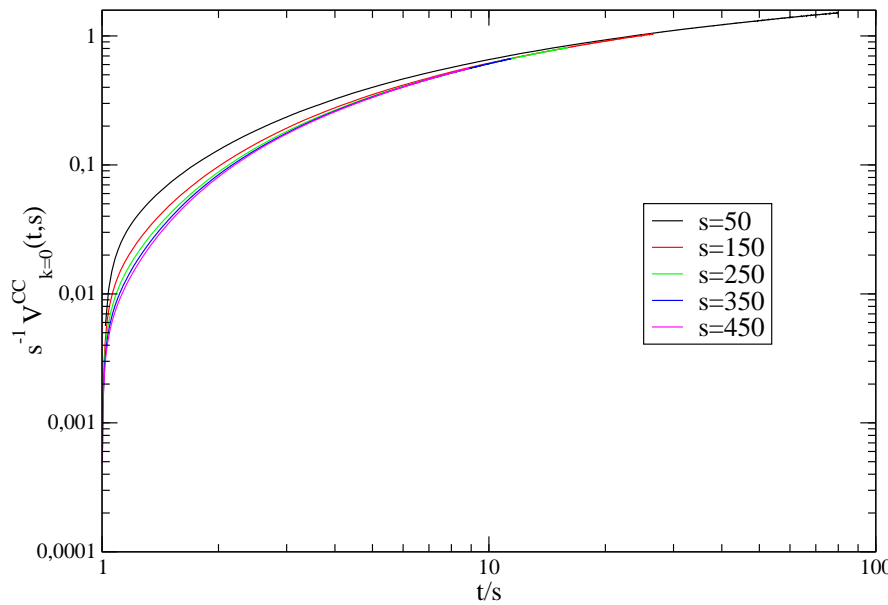


a

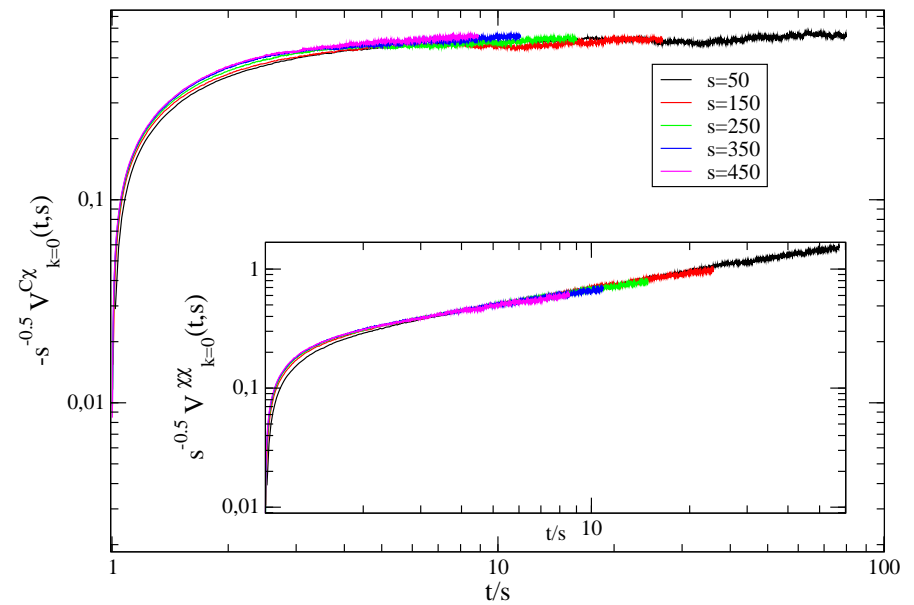


a F. Corberi, E. Lippiello, A. Sarracino, M. Zannetti, J. Stat. Mech. (2010) P04003

Numerical results Ising 2d



a



- Data are consistent with TRI being obeyed in 3d EA spin-glass but not in Ising model in $d = 2, 3$.

a F. Corberi, E. Lippiello, A. Sarracino, M. Zannetti, J. Stat. Mech. (2010) P04003

Restricted averages

- Given an average observable, like χ , writing $\chi = \sum_{\hat{C}} \chi_{\hat{C}} P(\hat{C})$, defines the *restricted average* $\chi_{\hat{C}}$.

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- For a variance it is:

$$V^{xx} = \sum_{\hat{C}} V_{\hat{C}}^{xx} P(\hat{C}) + \mathcal{V}^{xx},$$

where

$V_{\hat{C}}^{xx} = \langle \hat{\chi} \hat{\chi} \rangle_{\hat{C}} - \langle \hat{\chi} \rangle_{\hat{C}} \langle \hat{\chi} \rangle_{\hat{C}}$ (*restricted covariance* (\hat{C} fixed))

$\mathcal{V}^{xx} = \sum_{\hat{C}} \chi_{\hat{C}} \chi_{\hat{C}} P(\hat{C}) - \left(\sum_{\hat{C}} \chi_{\hat{C}} P(\hat{C}) \right)^2$
(covariance of the *restricted averages* $\chi_{\hat{C}}$ as \hat{C} is varied.)

Restricted averages and TRI

- Physical meaning: all sort of fluctuations of $\hat{\chi}_i$ are accounted for by the fluctuations of the local clock \hat{C}_i .

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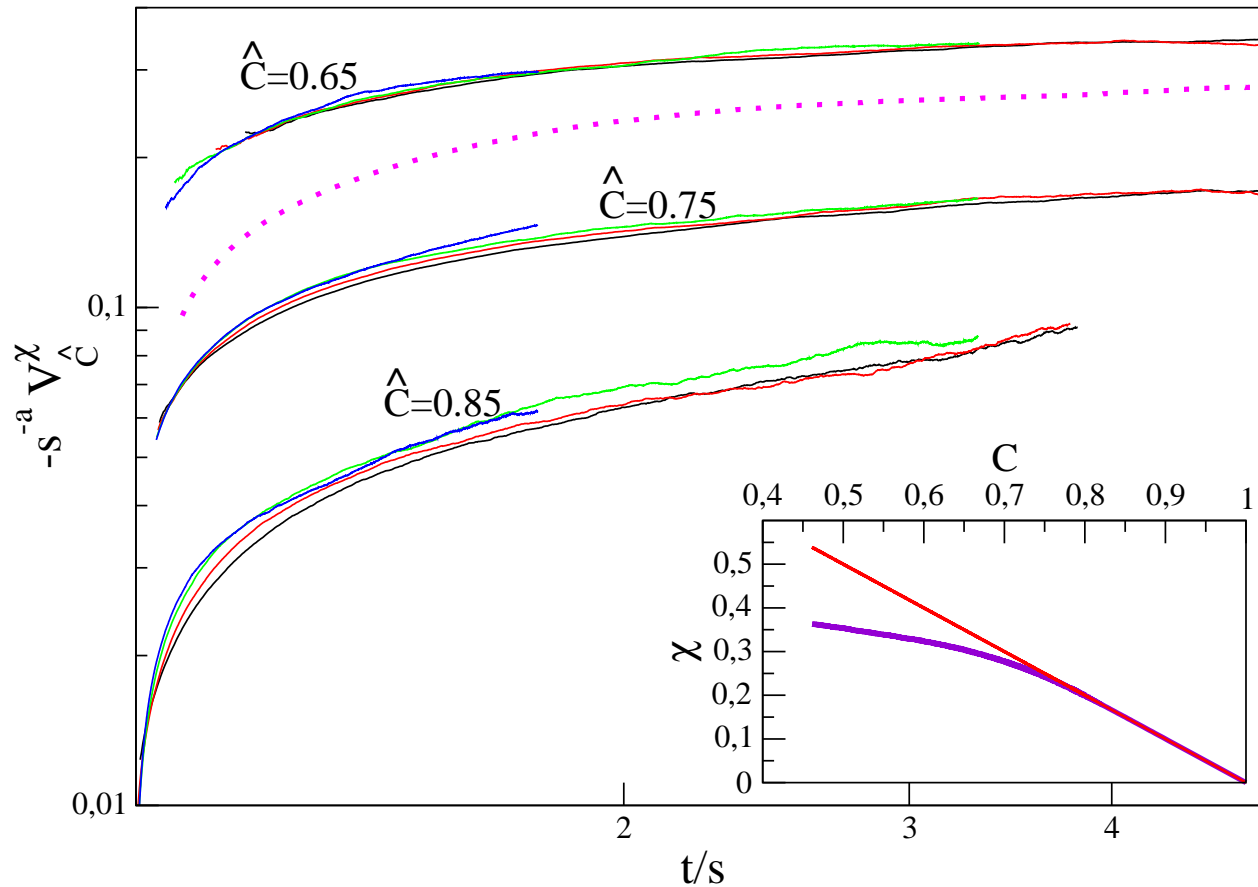
Restricted averages and TRI

- Physical meaning: all sort of fluctuations of $\hat{\chi}_i$ are accounted for by the fluctuations of the local clock \hat{C}_i .
- Hence one expects ^a

$$\sum_{\hat{C}} V_{\hat{C}}^{xx} P(\hat{C}) \ll \mathcal{V}^{xx},$$

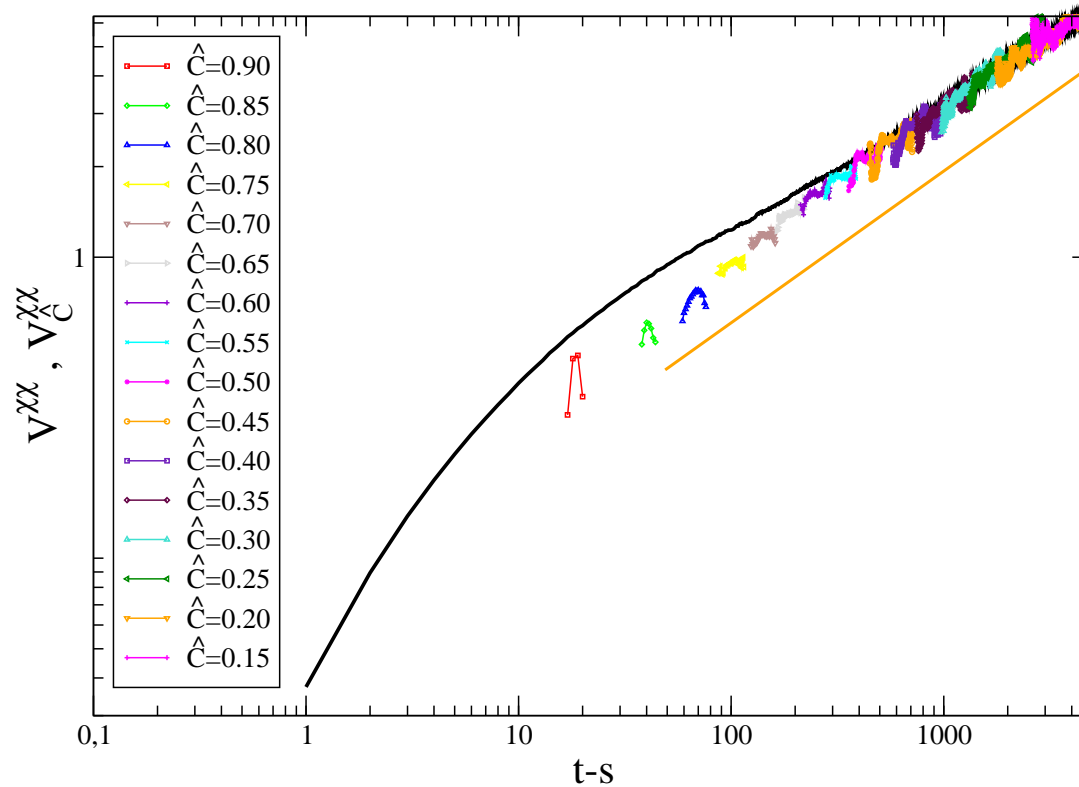
^a $V^{xx} = \sum_{\hat{C}} V_{\hat{C}}^{xx} P(\hat{C}) + \mathcal{V}^{xx}$

Restricted averages: numerics SG-3d



$$\sum_{\hat{C}} V_{\hat{C}}^{XY} P(\hat{C}) \ll \mathcal{V}^{XY}$$

Restricted averages: numerics Ising-2d



$$\sum_{\hat{C}} V_{\hat{C}}^{XY} P(\hat{C}) \gg \nu^{XY} \quad a$$

$$a \quad V^{XY} = \sum_{\hat{C}} V_{\hat{C}}^{XY} P(\hat{C}) + \nu^{XY}$$

Conclusions

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- Direct comparison between numerical simulations in EA SG and in Ising models shows that this is indeed the case. Even without resorting to TRI theory this is interesting in its own.
- Need for analytic results.

Thank you !